¹ Soft x-ray irradiation induced metallization of layered TiNCl

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- 16 Abstract

17 We have performed soft x-ray spectroscopy in order to study the photoirradiation time 18 dependence of the valence band structure and chemical states of layered transition metal 19 nitride chloride TiNCl. Under the soft x-ray irradiation, the intensities of the states near the Fermi level ($E_{\rm F}$) and the Ti³⁺ component increased, while the Cl 2p intensity decreased. 20 21 Ti 2p-3d resonance photoemission spectroscopy confirmed a distinctive Fermi edge with 22 Ti 3d character. These results indicate the photo-induced metallization originates from 23 deintercalation due to Cl desorption, and thus provide a new carrier doping method that 24 controls the conducting properties of TiNCl.

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26 Introduction

27 TiNCl is an interesting layered material that belongs to a family of layered nitride halides MNX(M = Ti, Zr, Hf and X = Cl, Br, I) systems with two polymorphs, the FeOCl-28 29 type crystal structure (α -form) and the SmSI-type crystal structure (β -form) [1, 2]. TiNCl 30 is known to take only the FeOCl-type crystal structure [2], consisting of a stack of 31 orthogonal M-N layers located between two X layers. TiNCl is a band insulator with a direct band gap of approximately 0.5-3 eV [3, 4], and theoretical studies have proposed 32 its application to optoelectronic devices [4], photocatalysts [5], and spin devices [6]. 33 34 Upon Na intercalation between TiNCl layers, TiNCl becomes a superconductor with a relatively high superconducting transition temperature of ~16 K [7]. As electron-doped 35

β-form HfNCl and ZrNCl [8-20], electron-doped TiNCl is considered a candidate for
 unconventional superconductors, where exotic mediation forces for Cooper pairing other

3 than phonon are discussed [21-25].

4 In order to further explore the physical properties of this remarkable material, it is 5 essential to control carrier concentrations. In the β -form MNX, the carrier control has 6 been performed by intercalation of an alkali- or alkaline-earth-metal [8, 9, 15, 16], and/or 7 off stoichiometry or deintercalation of X[26]. As for TiNCl, to the best of our knowledge, 8 carrier control by neither element substitution nor off stoichiometry has been reported. 9 This may be due to the fact that TiNCl is easily thermally decomposed into TiN by 10 annealing[27-29]. The intercalation of alkali metal atoms and/or organic molecules 11 between layers is the only method way for doping carriers. However, intercalated samples 12 are unstable in a humid air, and precise carrier controlling is challenging, which prevents 13 the systematic investigations of physical properties systematically.

14 One of the techniques to create conductive samples is photo-irradiation-induce 15 metallization by breaking chemical bonds [30-36]. In SrTiO₃ and TiO₂, a metallic state 16 called two-dimensional electron gases is created on the surface by photoirradiation. Two-17 dimensional electron gases have attracted extensive attention owing to unique physical 18 properties [30-35, 37-39], providing opportunities for the development of next generation 19 of electronic and photonic devices. If photo-induced bond breaking is also effective in 20 TiNCl, it can be established as a new method to control the physical properties of TiNCl. 21 Therefore, it is one of the important challenges to investigate whether TiNCl is metallized 22 by photoirradiation.

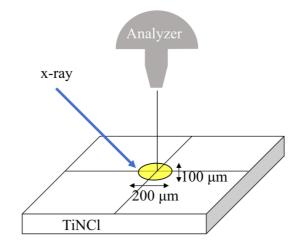
In this paper, we performed valence and core level photoelectron spectroscopy (PES), x-ray absorption spectroscopy (XAS), and Ti 2p-3d resonance photoelectron spectroscopy (RPES) of layered transition metal nitride chloride TiNCl. It was found that soft x-ray irradiation induced the metallization of TiNCl, and the metallization was closely related to the Cl desorption. This is the first report of electron-doped TiNCl due to Cl desorption.

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30 **Experiment**

The highly crystalline TiNCl was grown by a method described elsewhere [7]. It was pelletized into disk shape in an Ar filled glove box. Subsequently, the pellet was transferred to an ultra-high vacuum chamber for photoelectron spectroscopy without being exposed to the humid air for a long time. The experiments were performed at the soft x-ray beam line BL23SU of SPring-8 [40] using a photoemission spectrometer equipped with a Gammadata-Scienta SES-2002 electron analyzer. In the RPES 1 measurement, the energy resolution was set to 80 meV, and in the other PES 2 measurements, it was set to less than 200 meV. To reduce a possibility for pyrolysis of 3 TiNCl, we cooled down the sample and measured it at 100 K instead of at 300 K. We 4 performed PES measurements using a photon energy of 1000 eV. For calibration of the 5 binding energy, we used the Fermi edge of a gold film, which was located close to the 6 sample and had a good electrical contact to the sample.

The irradiation light of hv = 1000 eV was used, and the detector angle was set to 90°. 7 8 Sets of Ti 2p, N 1s, Cl 2p, and valence spectra were also measured at hv = 1000 eV in the 9 time interval of 25 minutes without changing the measured location of the sample surface. 10 The irradiation time includes the time of the PES measurements. Though it took 8 min to 11 take a set of Ti, N, Cl, and valence spectra, we refer to the irradiation time as the time 12 when we started the measurement of Ti. The spot size of the excitation light was 100 \times 200 μ m², and the photon flux was 2 × 10¹² Photons/sec [40]. The illustration of the 13 experimental geometry is shown in Fig. 1. The intensity of PES spectra was normalized 14 15 by the area intensity of Ti 2p core-level spectra, while those of RPES were normalized by 16 the photocurrent. XAS spectrum was measured using the total electron yield (TEY) mode. To obtain a clean surface, the TiNCl pellet was fractured at 100 K under an ultrahigh 17 vacuum of 1.5×10^{-8} Pa. 18



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Figure 1. Schematic diagram of the x-ray irradiation.

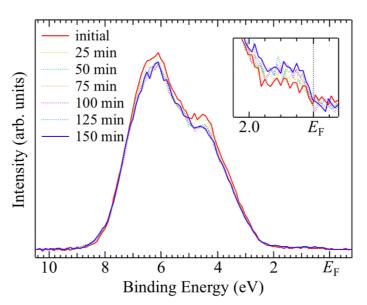
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22 **Results and discussion**

Figure 2 shows the soft x-ray irradiation time dependence of the valence band PES

1 spectra of TiNCl. The background of an iterative Shirley method [41, 42] was subtracted

- 2 from the raw spectra. The initial valence band spectrum of TiNCl had a peak at 6 eV with
- 3 a shoulder structure at 4 eV, and there was almost no intensity in the near- $E_{\rm F}$ region. The
- 4 overall spectral shape is consistent with that of the previous study [25], and thus the states
- 5 of approximately 6 and 4 eV can be ascribed to Cl 3*p* and N 2*p* orbitals, respectively.
- The peak intensities at 6 eV and 4 eV were decreased gradually. In contrast, the 6 7 spectral intensity in the region between $E_{\rm F}$ and ~ 1.5 eV appears to be increased, as 8 illustrated in the inset that shows a blowup of the spectra near $E_{\rm F}$. The intensity near $E_{\rm F}$ 9 of the initial spectrum can be ascribed to the states induced by the soft x-ray irradiation 10 during the measurement and/or may be assigned to impurity states. In the spectrum after 11 150 min irradiation, a clear Fermi edge structure was observed, suggesting the metallic nature of the measured region of the sample surface. The change in the valence band and 12 13 the intensity near $E_{\rm F}$ due to the soft x-ray irradiation is very similar to the change due to the photoirradiation effect of SrTiO₃ [30]. 14



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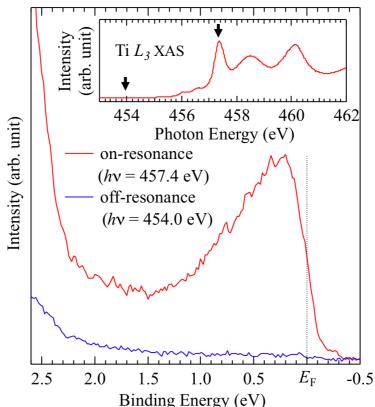
16 Figure 2. Soft x-ray irradiation time dependence of the valence band PES spectra of TiNCl,

17 measured with a photon energy of 1000 eV at 100 K. Inset shows the enlargement of the 18 spectra near $E_{\rm F}$.

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To spectroscopically confirm the metallization and the character of the states at E_F , we performed RPES of irradiated TiNCl at the Ti 2p-3*d* absorption threshold. The photon energies for the RPES measurement were chosen based on the Ti L_3 XAS spectrum, as shown in the inset of Fig. 3. Ti L_3 XAS spectrum had sharp absorption peaks at 457.4 eV, 458.5 eV, and 460.1 eV in the region of Ti $2p_{3/2}$ -3*d* absorption edge. In the octahedral

1 structure, the 3d level was split into e_g and t_{2g} states by the crystal field effect. However, 2 since the Ti ions in TiNCl are in a distorted octahedral structure [7], the degeneracy was 3 further removed by the static Jahn-Teller effect, leading to the observed three-peak structure. We measured RPES spectra of TiNCl in the off-resonance (hv = 454.0 eV) and 4 5 on-resonance (hv = 457.4 eV) conditions after the irradiation the soft x-rays for 4 h. While 6 the spectral intensity at $E_{\rm F}$ in the off-resonance spectrum was almost negligible, the one 7 in the on-resonance spectrum exhibited a distinctive Fermi edge, and was considerably 8 enhanced. The resonance enhancement of the near- $E_{\rm F}$ structure indicates that its orbital 9 character is Ti 3d. Furthermore, the presence of the distinctive Fermi edge indicates that 10 the soft x-ray irradiation indeed induces the metallization of TiNCl.



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Figure 3. ON (red) and OFF (blue) Ti 2p-3d resonance photoemission spectra of the irradiated TiNCl measured with a photon energy of 457.4 eV and 454.0 eV, respectively. The inset shows the Ti 2p absorption spectrum where the employed photon energies for the resonance photoemission are indicated by arrows.

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To understand the mechanism of the metallization induced by soft x-ray irradiation in terms of chemical states, we have measured the core level spectra of TiNCl with different irradiation conditions. Figure 4 (a) shows soft x-ray irradiation time dependence of the Ti 20 2p core-level spectra, and Fig.4 (b) represents the initial and 150 min difference spectrum.

The peaks at 457.5 eV and 463.5 eV are spin-orbit split Ti 2p, namely Ti $2p_{3/2}$ and Ti $2p_{1/2}$, 1 2 respectively [43]. The initial Ti 2p spectrum showed a sharp Ti⁴⁺ peak at 457.5 eV [25] and a small shoulder structure at 456.5 eV. In the Ti 2p spectrum, as the irradiation time 3 increased, the peak area of Ti⁴⁺ decreased compared to the initial value. In addition, the 4 intensity of the shoulder structure at 456.5 eV clearly increased gradually. This binding 5 energy is different from that of TiN ($455.2 \pm 0.2 \text{ eV}$) [44, 45], but is very similar to the 6 7 week Ti 2p spectral component of electron-doped SrTiO₃ and electron-doped anatase-TiO₂, which is assigned to the component of Ti³⁺ state [30, 31, 46-48]. Thus, the 8 observation indicates that the soft x-ray irradiation on TiNCl induces the changes Ti⁴⁺ in 9 the valence state of the Ti ion from the Ti³⁺ state to the Ti⁴⁺ state. We performed spectral 10 fitting using two components. The area intensities of Ti³⁺ and Ti⁴⁺ components at each 11 irradiation time are summarized in Table 1. 12

13 Figure 4 (c) compares the spectra of the N 1s core level at each soft x-ray irradiation time. These spectra had peaks at 396.5 eV, and were almost identical, suggesting that 14 15 there was no significant change in the spectral shape due to soft x-ray irradiation. The Cl 2p core-level spectra consist of Cl $2p_{3/2}$ (198.4 eV) and Cl $2p_{1/2}$ (200.1 eV) spin-orbit split 16 peaks. The intensity of the Cl 2p spectrum after the soft x-ray irradiation of 150 min 17 decreased gradually, which suggests the desorption of Cl atoms from the surface is 18 19 induced by the soft x-ray irradiation. This is consistent with the observed reduction of the 20 valence band spectra, which reflects a dominant Cl 3p contribution.

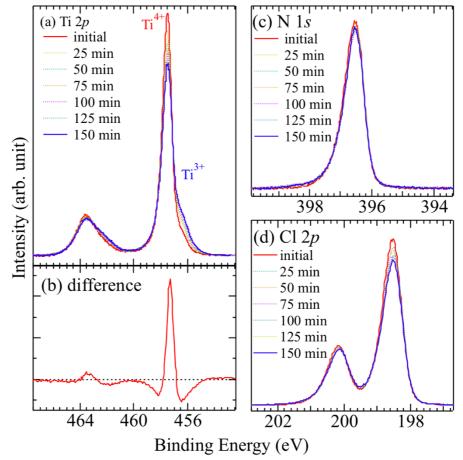
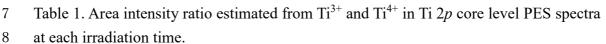


Figure 4. Soft x-ray irradiation time dependence of core-levels PES spectra. (a) Ti 2p
core-level spectra irradiated from the initial to 150 min in the interval of 25 min. (b)
Difference between the initial and 150min. (c) and (d) are analogous spectra to (a) for the

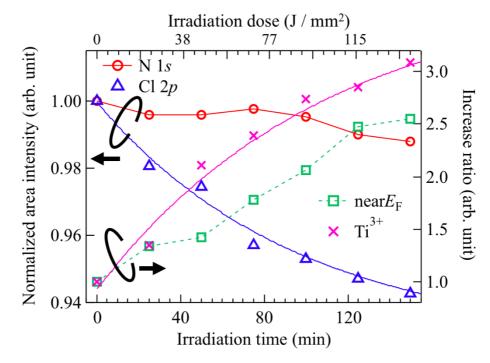
- 5 N 1s and Cl 2p core-level, respectively.



Irradiation time (min)	Relative area intensity of components (%)	
	${ m Ti}^{4+}$ in Ti $2p$	${\rm Ti}^{3+}$ in Ti $2p$
0	90.60	9.40
25	87.37	12.63
50	80.22	19.78
75	77.52	22.48
100	74.26	25.74
125	73.21	26.79
150	70.96	29.04

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2 Here, we discuss the relationship between the observed results. Figure 5 shows the soft x-ray irradiation time dependence of the intensity of N 1s and Cl 2p core-level PES 3 spectra, the PES intensity between $E_{\rm F}$ and 1.5 eV, and the Ti³⁺ ratio (Ti³⁺ / (Ti⁴⁺ + Ti³⁺)) 4 of Ti 2p core-level PES spectrum. We normalized their initial values to unity. The 5 6 intensity of N 1s core-level PES does not show a marked change depending on the soft 7 x-ray irradiation time of soft x-ray. In a comparison with that of the N 1s core-level PES, the intensity of the Cl 2p core-level PES decreased more rapidly with increasing 8 irradiation time. Furthermore, simultaneously with a decrease in Cl, the areas of the Ti³⁺ 9 component and of the states near $E_{\rm F}$ increased in a similar manner. A smaller increase in 10 near $E_{\rm F}$ intensity may be attributed to a more complicated orbital character of states near 11 $E_{\rm F}$ than that of the Ti³⁺ core level. These observations suggest the close relationship of 12 the desorption of Cl atoms and an increase in the intensity near $E_{\rm F}$ with an increase in the 13 Ti³⁺ component. We consider that the desorption of a portion of the Cl atoms from the 14 TiNCl layers introduces electron carriers into the layers, making the layer conductive, 15 which is inferred from the observation of the metallic Fermi edge. Therefore, to estimate 16 the time of change, the data for Ti^{3+} and Cl 2p, which are more intense than the near- E_F 17 data, were fitted with an exponential function $(A \exp(-t/\lambda) + B \text{ for } \text{Ti}^{3+} \text{ and } C \exp(t/\lambda) + D \text{ for } \text{ and } C \exp(t/\lambda) + D \text{ for } \text{ and } C \exp(t/\lambda) + D \text{ for } \text{ and } C \exp(t/\lambda) + D \text{ for } \text{ for } C \exp(t/\lambda) + D \text{ for } \text{ for } C \exp(t/\lambda) + D \text{ for } \text{ for } C \exp(t/\lambda) + D \text{ for } C \exp(t/\lambda)$ 18 Cl 2p, where t and λ are time and lifetime, respectively.). The obtained lifetimes of Ti³⁺ 19 and Cl 2p were 119 \pm 29 and 96 \pm 31 min, respectively. This provided a value of 100 20 min as the order estimate of the time of change for this experiment. 21



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1 Figure 5. Soft x-ray irradiation time dependence of the core-level and near- $E_{\rm F}$ intensities

- of TiNCl. Area intensity of N 1s and Cl 2p core-level spectra (left axis), and an increase ratio in area intensity of near E_F and Ti³⁺ in Ti 2p core level spectra (right axis). The initial values were normalized to 1.
- 5

6 There are two possible causes of the desorption of Cl atoms. One of them is a 7 pyrolysis. TiNCl has been reported to be decomposed completely into TiN when it is 8 heated above 550°C [27-29]. In the present experiment, however, the sample was kept at 9 100 K during the measurement, and therefore this possibility can be ruled out. The other 10 possible cause is photon stimulated desorption [49-51]. There are primarily two 11 mechanisms for photon stimulated desorption, which are distinguished by the process 12 induced by photon irradiation: The Menzel, Gomer, and Redhead (MGR) model and the 13 Knotek and Feibelman (KF) model [49, 50]. The MGR model involves the excitation of 14 valence electrons and works reasonably well for covalent bonds. The KF model involves 15 the excitation of core electrons and is applicable to ion bonding. Since the TiNCl is 16 considered to be an ionic insulator composed of TiN⁺ and Cl⁻ [23], we assumed that this mechanism is involved in the Cl desorption of TiNCl. In the case of TiO₂, photoinduced 17 desorption is explained by the KF model in terms of Coulomb expulsion between Ti⁴⁺ 18 19 and O⁺ that are induced by emitting Auger electrons after the relaxation of electrons in O 20 2p levels to create photo-holes in the Ti 3p levels [48]. If the KF mechanism is involved 21 in the present case, the electron transfer that leads to Coulomb expulsion may occur from 22 Cl 3p (valence band) to Ti 3p. Further experiments are needed to confirm the KF 23 mechanism, as it is we think beyond the scope of the present study.

24 With photon stimulated Cl desorption, the final products may be TiN or Cl deficient 25 TiNCl₁₋₆, depending on the amount of Cl deficiency. However, from the Ti 2p 26 photoelectron spectrum, no sharp TiN peak was observed at 455±0.3 eV[44, 45], 27 suggesting the formation of TiNCl_{1- δ}, rather than TiN. In β -(Hf or Zr) NCl, Cl deficiency is called deintercalation, which is one of the methods of electron doping to the system. 28 29 Deintercalated HfNCl_{0.7} is a metallic system that also exhibits superconductivity below 24 K [26]. However, since electron doping utilizing the Cl deintercalation into α -MNX 30 31 has never been reported, the present study is the first experimental realization of 32 deintercalated electron doped α -MNX. Further Cl deintercalation by the soft x-ray 33 irradiation would make the system superconductive, like β -form compounds.

Lastly, we discuss the benefits of methodologies and findings. Regarding pyrolysis, bond breaking by heat cannot be confined within a small area because of the diffusion of heat in a solid. In addition, temperature controlling of a sample is severe for a target

material located in a narrow temperature region of a complicated phase diagram. In 1 2 contrast, bond breaking by light irradiation occurs only in the place where light is 3 irradiated. In particular, when core-levels excitation are important in the bond-breaking process as in the KF model, element/chemical site-selective bond breaking is possible by 4 5 setting the excitation energy to a specific element/chemical site [52, 53]. This site-6 selective bond breaking is called a molecular scalpel and is mainly studied in surface 7 adsorption molecular systems. Such a characteristic, combined with sub-micrometer scale 8 spot sizes in third-generation synchrotron facilities, might be used to make sub-micron 9 scale conducting paths at any place on the surface. Photo-induced metallization of TiNCl 10 has enabled detailed electronic-structure investigation studies of metallic (even 11 superconductive) TiNCl by controlling the carrier concentration in a small single crystal 12 region on the surface, which shed light on the exotic properties of TiNCl.

13

14 **Conclusion**

15 In summary, the irradiation time dependence of the valence band and core-level spectra of TiNCl was studied by PES and Ti 2p-3d RPES. Using soft x-ray irradiation, 16 17 the intensity of the valence band structure origination from Cl 3p and N 2p states decreased while the intensity near E_F increased. The intensity near E_F exhibited resonant 18 19 enhancement, and the Fermi edge was clearly observed by Ti 2p-3d RPES. The analyses 20 of core-level and near- E_F PES spectral intensities as a function of the soft x-ray irradiation 21 time revealed that the desorption of Cl atoms, the increase in the intensity of the Ti³⁺ 22 component, and increase of the spectral weight in the vicinity of $E_{\rm F}$ exhibit strong 23 correlations. These results indicate that photoirradiation induced metallization of TiNCl occurs due to electron doping through Cl desorption. As estimated from Cl 2p and Ti³⁺ 24 25 data, the order of the time-of-changes is 100 min. This technique can be used to further 26 explore the conducting properties, especially the unconventional superconductivity, of 27 TiNCl.

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20 21 22 23 24 25 26 27 28 29 30 31 32 33 34	 38. 39. 40. 41. 42. 	117602 Santander-Syro A <i>et al</i> 2011 Two-dimensional electron gas with universal subbands at the surface of SrTiO ₃ 469 189 Santander-Syro A <i>et al</i> 2014 Giant spin splitting of the two-dimensional electron gas at the surface of SrTiO ₃ <i>Nature Nat. Mater.</i> 13 1085 King P <i>et al</i> 2014 Quasiparticle dynamics and spin–orbital texture of the SrTiO ₃ two- dimensional electron gas <i>Nat. Commun.</i> 5 3414 Saitoh Y <i>et al</i> 2012 Performance upgrade in the JAEA actinide science beamline BL23SU at SPring-8 with a new twin-helical undulator <i>J. Synchrotron Rad.</i> 19 388 Shirley D 1972 High-Resolution X-Ray Photoemission Spectrum of the Valence Bands of Gold <i>Phys. Rev.</i> B 5 4709 Proctor A, Sherwood P 1982 Data Analysis Techniques in X-ray Photoelectron Spectroscopy <i>Anal. Chem.</i> 54 13 The binding energies of Ti 2 <i>p</i> and Cl 2 <i>p</i> of the present study are slightly lower by 0.4 eV and 0.2 eV than those in Ref. 25, respectively. The two experiments were

charging up, difference in metallization, and calibration of the spectrometer. Since
the photon flux density of BL25SU is higher than that of BL23SU, larger binding
energies of BL25SU data may be explained by charging up and/or difference in
metallization. We do not have the clear answer, however the difference does not
change the conclusion of the present study.
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