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Fabrication of field-effect transistor device with higher fullerene, C<sub>88</sub>

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A fullerene field-effect transistor (FET) device has been fabricated with thin films of  $C_{88}$ , and n-channel normally-on depletion-type FET properties have been found in this FET device. The  $C_{88}$  FET exhibited a high mobility,  $\mu$ , of 2.5 x 10<sup>-3</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> at 300 K, in fullerene FETs. The carrier transport showed a thermally-activated hopping transport. The n-channel normally-on FET properties and the hopping transport reflect the small mobility gap and low carrier concentration in the channel region of  $C_{88}$  thin-films.

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A number of field-effect transistors (FETs) have been fabricated with thin films of organic molecules, and their properties have been studied for the next-generation electronics during the past 10 years.<sup>1,2</sup> The mobilities,  $\mu$ , of the FETs with organic molecules (organic FETs) are generally lower than those, ~ 1000 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, of the conventional FETs with inorganic materials (inorganic FETs). However, the organic FETs are well known to have many advantages such as large-area coverage, structural flexibility, and low-temperature and low-cost processing, in comparison with inorganic FETs.

The highest  $\mu$  value in p-channel organic thin-film FET devices was 1.5 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for the pentacene thin-film FET.<sup>3</sup> Recently, the  $\mu$  value reached to 8 - 20 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> in the rubrene single-crystal FET device exhibiting p-channel behavior.<sup>4-6</sup> On the other hand, the highest  $\mu$  value in n-channel organic FET devices was 0.6 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for the *N*,*N*'-dialkyl-3,4,9,10-perylene tetracalboxylic diimide derivative FET device.<sup>7</sup> These  $\mu$  values are comparable to those of the recent amorphous Si FETs.<sup>1,2</sup>

The fullerene FET was first fabricated with thin films of  $C_{60}$  by Haddon *et al.*<sup>8</sup> This FET device showed the high  $\mu$  value of  $0.08 - 0.3 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ . Subsequently, the  $C_{70}$  FET device was fabricated and the n-channel properties were observed with the  $\mu$  of 2 x  $10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1.9}$  Recently, the  $\mu$  value for the  $C_{60}$  thin-film FET reached to  $0.56 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  when the FET properties were measured under  $10^{-9}$  Torr without an exposure to air.<sup>10</sup> Thus, the high  $\mu$  value implies that the  $C_{60}$  FET can play an important role in future applications towards the practical organic FET devices.

We have examined to fabricate the FET devices with thin films of various types of higher fullerenes and endohedral metallofullerenes in order to improve their FET properties.<sup>11-14</sup> All

FET devices with higher fullerenes and endohedral metallofullerenes showed n-channel normally-on properties which originated from high bulk currents caused by small gap energies,  $E_g$ 's. Thus, the observed properties reflect the intrinsic nature of fullerenes used as active layer. Therefore, the fabrications and characterizations of the FET devices with a series of higher fullerenes and endohedral metallofullerenes are indispensable to know their electronic properties and possibility of application as electronic materials. In the present study, we have first fabricated a high-performance FET device with higher fullerene, C<sub>88</sub>. The electronic properties of giant fullerenes such as C<sub>88</sub> have never been studied because of a difficulty in obtaining pure solid-samples. Therefore, the present study is the first approach to the electronic properties of giant fullerenes.

The schematic representation of  $C_2$ -isomer (#7),  $C_2(7)$ , of  $C_{88}$  and the cross-sectional view of the  $C_{88}$  FET device are shown in Fig. 1(a); the  $C_2(7)$ - $C_{88}$  is one of seven  $C_2$  isomers of  $C_{88}$ .<sup>15</sup> The  $C_{88}$  sample was purified by high performance liquid chromatography with toluene as eluent, after a Soxhlet-extraction of the soot containing  $C_{88}$  with 1,2,4-*tri*-chlorobenzene, where the soot was prepared by an arc-discharge of carbon graphite rods containing Eu<sub>2</sub>O<sub>3</sub>. The time-of-flight mass spectrum of the purified sample is shown in Fig. 1(b), which exhibited only a single peak for  $C_{88}$ . The purified sample may contain some isomers since the isomer-separation of  $C_{88}$  sample was not performed. The commercially available SiO<sub>2</sub>/Si(100) wafer was used as a substrate after washing with acetone, methanol and H<sub>2</sub>SO<sub>4</sub>/H<sub>2</sub>O<sub>2</sub>. The SiO<sub>2</sub>/Si substrate was treated with hexamethyldisilazane (HMDS) to form a hydrophobic surface. The thickness and capacitance,  $C_0$ , of SiO<sub>2</sub> were 420 nm and 8.2 x 10<sup>-9</sup> F cm<sup>-2</sup>, respectively. The thin film of C<sub>88</sub> was formed by a thermal deposition under a vacuum of 10<sup>-8</sup> Torr. The channel length, *L*, and the channel width, *W*, of this device were 30 and 3000  $\mu$ m, respectively. The FET properties of the C<sub>88</sub> FET device were measured under 10<sup>-6</sup> Torr after annealing for 13 h at 120°C under 10<sup>-6</sup> Torr.

The drain current,  $I_D$ , *vs.* drain-source voltage,  $V_{DS}$ , plots for the C<sub>88</sub> FET at 300 K are shown in Fig. 1(c). The plots show n-channel normally-on depletion-type FET properties. The  $I_D - V_G$  plot at  $V_{DS} = 20$  V is also shown in Fig. 1(d). The  $I_D$  increases with an increase in  $V_G$  from -30 to 100 V. It should be noted that the channel conductance decreases by applying the negative  $V_G$ , due to the depletion of carriers in the channel. The  $\mu$  value at 300 K for the C<sub>88</sub> FET was estimated to be 2.5 x 10<sup>-3</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> from the  $I_D$  -  $V_G$  plot above  $V_G = 60$  V (Fig. 1(d)) with the equation (1) for  $V_{DS} \ll V_G - V_T$ ,

$$I_{\rm D} \sim (\mu W C_0 / L) (V_{\rm G} - V_{\rm T}) V_{\rm DS}$$
(1)

where  $V_{\rm T}$  refers to the threshold voltage.<sup>16</sup> This  $\mu$  value is the highest one next to that of the C<sub>60</sub> FET among fullerene FETs. The  $V_{\rm T}$  was estimated to be –142 V from the  $I_{\rm D}$ - $V_{\rm G}$  plot (Fig. 1(d)). Furthermore, the  $\mu$  and  $V_{\rm T}$  values were estimated from the  $I_{\rm D}$  –  $V_{\rm G}$  plots at  $V_{\rm DS}$  of 5, 10 and 30 V in order to confirm the validity of the very high negative value of  $V_{\rm T}$ . The  $\mu$  and  $V_{\rm T}$  values were almost consistent with each other, and the same as those at  $V_{\rm DS} = 20$  V, showing that the  $\mu$  and  $V_{\rm T}$  values are reliable. The negative  $V_{\rm T}$  of -142 V supports that the C<sub>88</sub> FET is normally-on type. Such a depletion-type property was observed for C<sub>82</sub> and C<sub>84</sub> FETs.<sup>12,13</sup> The on-off ratio was estimated to be ~7 from the  $I_{\rm D}$  at  $V_{\rm G} = 100$  V and that at  $V_{\rm G} = -30$  V. This value is also the best one at room temperature among normally-on type FET devices: ~2.1 for C<sub>82</sub>, ~6 for C<sub>84</sub> and ~3.1 for Pr@C<sub>82</sub>.<sup>12-14</sup> These FET parameters show that C<sub>88</sub> is the second

best material for active layer of FET among fullerenes.

Here, we have investigated the variation of  $\mu$  for the applied V<sub>G</sub> with the equation (2),

$$\mu \sim (L / (WC_0 V_{\rm DS})) \, \mathrm{d}I_{\rm D} / \mathrm{d}V_{\rm G} \tag{2}$$

in order to confirm whether the FET parameters obtained reflect the intrinsic carrier transport in the C<sub>88</sub> FET; if the  $\mu$  value is not influenced by the V<sub>G</sub>, the FET parameters, inclusive of the  $\mu$ , can closely be associated with the intrinsic nature of the carrier transport in the C<sub>88</sub> FET device.<sup>17</sup> The V<sub>G</sub> dependence of  $\mu$  value estimated with eq. (2) is shown in Fig. 2, and the  $\mu$ value above V<sub>G</sub> of 60 V is almost the same as the  $\mu$ , 2.5 x 10<sup>-3</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, estimated with eq. (1). The  $\mu$  value ranges from 2 x 10<sup>-3</sup> to 6.5 x 10<sup>-3</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> at V<sub>G</sub> of -30 ~ 100 V, as shown in Fig. 2. This result implies that the FET properties are partly influenced by the V<sub>G</sub>-dependence of resistance of source and drain contacts owing to two-probe measurements used in this study.

The  $I_D - V_{DS}$  plots at 230 K for the C<sub>88</sub> FET is shown in Fig. 3(a), and the  $I_D - V_G$  plot is shown in Fig. 3(b). The  $\mu$  and  $V_T$  were estimated to be 7.2 x 10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> and -50 V, respectively, with eq. (1). The on-off ratio was 24, whose value was larger by three times than that at 300 K. The  $I_D$  at  $V_G = 0$  V was lower by one order of magnitude than that at 300 K. Thus a decrease in bulk current can lead to an increase in the on-off ratio at 230 K. The  $\mu$ increased exponentially with an increase in temperature, *T*, up to 350 K (Fig. 3(c)), and the  $\mu$ value reached to 8.0 x 10<sup>-3</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> at 350 K. The plot of  $\mu$  - *T* shows that the channel conduction of the C<sub>88</sub> FET device follows a thermally-activated hopping-transport model ( $\mu$  ~ exp- $E_a/k_BT$ ). The activation energy  $E_a$  was estimated to be 0.13 eV from the ln $\mu$  - 1/*T* plot shown in the inset of Fig. 3(c). This value is consistent with those, 0.14 and 0.13 eV, for the C<sub>82</sub> and C<sub>84</sub> FETs, respectively,<sup>12,13</sup> while the value is smaller than that, 0.22 eV, for the Pr@C<sub>82</sub> FET.<sup>14</sup> These results show that the FET properties of the higher fullerene FETs are similar with each other. Further, we can point out that the properties of higher fullerenes are different from those of endohedral metallofullerenes; for example, the  $\mu$  values, ~ 10<sup>-3</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, of higher fullerene FETs are higher by one order of magnitude than those, ~10<sup>-4</sup> cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, of endohedral metallofullerenes, and the depletion properties can be more clearly observed in the higher fullerenes than endohedral metallofullerenes.<sup>11-14</sup> The high  $\mu$  for the higher fullerenes originates from the fact that intermolecular overlap of the lowest unoccupied molecular orbitals (LUMOs) of the higher fullerenes is large because of the delocalized  $\pi$  character. Since more critical factor is the crystallinity of the thin films, it is necessary to investigate the crystallinity of thin films of higher fullerenes and endohedral metallofullerenes in detail.

The bulk current increases exponentially with increasing *T*, where the  $I_D$  at  $V_G = 0$  V and  $V_{DS} = 20$  V was used as bulk current. The gap energy,  $E_g$ , was estimated to be 0.42 eV from the  $\ln I_D - 1/T$  plot, and the  $E_g$  is consistent with those, 0.43 and 0.55 eV, for C<sub>82</sub> and C<sub>84</sub>, respectively, estimated from resistivity  $\rho - T$  plot.<sup>12,13</sup> The  $E_g$  cannot be attributed to the band gap but mobility gap, as suggested by Shiraishi et al.<sup>18</sup> The  $V_T$  decreases monotonically with increasing *T* up to 350 K (Fig. 3(d)). Such *T* dependence can be reasonably explained by an increase in bulk current described above. The  $V_T$  increased to ~0 V at 160 K. Thus, the FET property changes from normally-on to normally-off-like with decreasing *T* owing to a

decrease in bulk current.

In the present study, some intrinsic characters for the  $C_{88}$  FET have been clarified, although the  $V_G$  dependence of resistance of source and drain contracts partly affects the FET properties observed. The normally-on FET properties found in this device reflects the small  $E_g$  for the  $C_{88}$  thin-films. The delocalized LUMOs in  $C_{88}$  produce the larger  $\mu$  than those in the endohedral metallofullerene FETs. Furthermore, the hopping conduction observed in the  $C_{88}$  thin-films originates from low carrier concentration of at most 5 x  $10^{12}$  cm<sup>-2</sup> at  $V_G = 100$  V, *i.e.*, the metallic conduction never occurs at such low concentration because 0.05 electrons are exclusively injected into a  $C_{88}$  molecule. Therefore, the high carrier injection is very important for a realization of metallic transition, and one of the approaches to high carrier injection is to use high dielectric materials for the FET devices. In the present study, the electronic properties of  $C_{88}$  have been investigated, and its possibility of applications for electronic materials has been indicated from the highest FET performance next to the  $C_{60}$  FET.

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Figure captions

Fig. 1. (a) Schematic representation of  $C_{88}(C_2(7))$  and cross sectional view of  $C_{88}$  FET device. (b) Time-of-flight (TOF) mass spectrum of the purified sample. (c)  $I_D$ - $V_{DS}$  and (d)  $I_D$ - $V_G$  plots at  $V_{DS} = 20$  V for  $C_{88}$  FET at 300 K. Closed circles refer to the points measured at  $V_G = -30 \sim 100$  V. Fig. 2.  $V_{\rm G}$  dependence of  $\mu$  for the C<sub>88</sub> FET at 300 K.

Fig. 3. (a)  $I_D$ - $V_{DS}$  and (b)  $I_D$  -  $V_G$  plots at  $V_{DS} = 20$  V for C<sub>88</sub> FET at 230 K. Closed circles refer to the points measured at  $V_G = -30 \sim 100$  V. Plots of (c)  $\mu$  - T and (d)  $V_T$  - T. ln  $\mu$  - 1/T plot in the inset of (c).

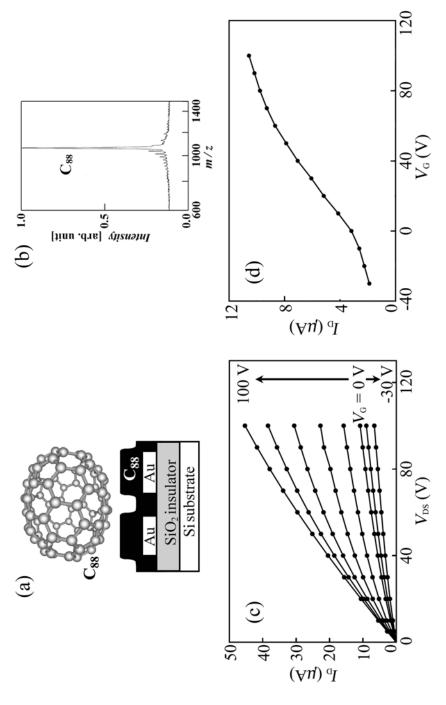


Figure 1.

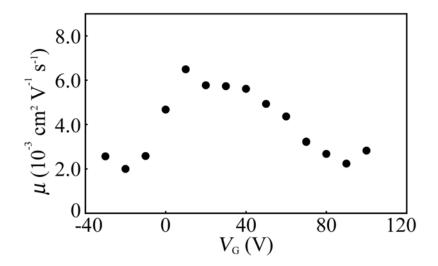


Figure. 2

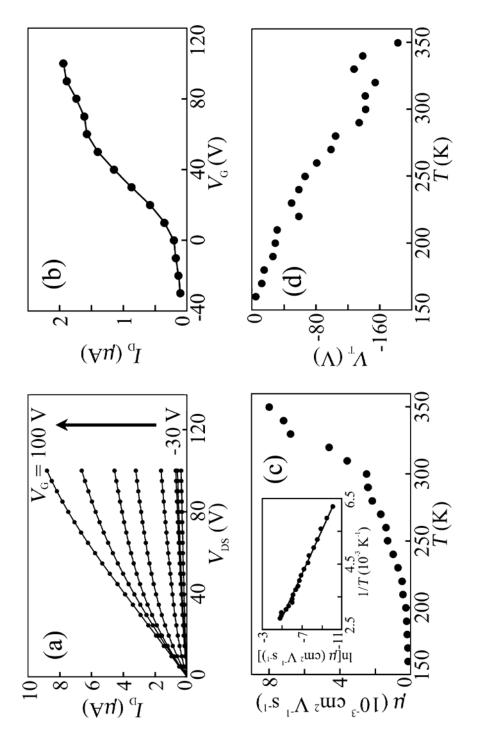


Figure. 3